2016-2 QUANTUM MECHANICS FOR ELECTRICAL AND ELECTRONIC ENGINEERS

# **Wave Propagation**

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### Background

a. Plane waves and interference

#### **b.** Diffraction

c. Diffraction from periodic structures

#### Summary





## Background

#### In the previous lecture...

#### 1-D wave equation

$$\frac{\partial^2 y}{\partial z^2} - \frac{1}{c^2} \frac{\partial^2 y}{\partial t^2} = 0 \quad \text{c: wave velocity}$$



For monochromatic wave (single angular frequency  $\omega$ )

 $\phi(z,t) = Z(z)T(t)$ 

$$\frac{\partial^2 \phi(z,t)}{\partial t^2} = -\omega^2 \phi(z,t)$$

 $\frac{d^2 Z(z)}{dz^2} + k^2 Z(z) = 0 \qquad k^2 = \frac{\omega^2}{c^2} \qquad : \text{Helmholtz wave equation}$ 

ω:angular frequency





## a. Plane waves and interference

#### Wave equation in 3 dimensions

$$\nabla^2 \phi(x, y, z, t) - \frac{1}{c^2} \frac{\partial^2 \phi(x, y, z, t)}{\partial t^2} = 0$$
  
where  $\nabla^2 \equiv \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$ 

### For monochromatic wave (single angular frequency $\omega$ )

Verifying plane wave solutions of the form when k=  $\omega$  /c

$$\phi(x, y, z, t) = \exp[i(\mathbf{k} \cdot \mathbf{r} - \omega t)]$$

where 
$$\mathbf{r} = x\,\hat{\mathbf{x}} + y\,\hat{\mathbf{y}} + z\,\hat{\mathbf{z}}$$
  $\mathbf{k} = k_x\,\hat{\mathbf{x}} + k_y\,\hat{\mathbf{y}} + k_z\,\hat{\mathbf{z}}$ 



#### (1955) (1156/5196) (1210-9) (1210-9)

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## a. Plane waves and interference

#### For monochromatic wave (single angular frequency $\omega$ )

Verifying plane wave solutions of the form when k=  $\omega$  /c

$$\nabla \exp[i(\mathbf{k} \cdot \mathbf{r} - \omega t)] = \left(\hat{\mathbf{x}}\frac{\partial}{\partial x} + \hat{\mathbf{y}}\frac{\partial}{\partial y} + \hat{\mathbf{z}}\frac{\partial}{\partial z}\right) \exp\left\{i\left[k_x x + k_y y + k_z z - \omega t\right]\right\}$$
$$= i\left(k_x \hat{\mathbf{x}} + k_y \hat{\mathbf{y}} + k_z \hat{\mathbf{z}}\right) \exp\left\{i\left[k_x x + k_y y + k_z z - \omega t\right]\right\}$$

$$= i\mathbf{k} \exp[i(\mathbf{k} \cdot \mathbf{r} - \omega t)]$$

$$\nabla^{2} \exp[i(\mathbf{k} \cdot \mathbf{r} - \omega t)] = \nabla \cdot \nabla \exp[i(\mathbf{k} \cdot \mathbf{r} - \omega t)]$$

$$= \nabla \cdot \left(i\vec{k} \exp[i(\mathbf{k} \cdot \mathbf{r} - \omega t)]\right)$$

$$= i\left(k_{x}\frac{\partial}{\partial x} + k_{y}\frac{\partial}{\partial x} + k_{z}\frac{\partial}{\partial x}\right)\exp[i(\mathbf{k} \cdot \mathbf{r} - \omega t)]$$

$$= -(k_{x}^{2} + k_{y}^{2} + k_{z}^{2})\exp[i(\mathbf{k} \cdot \mathbf{r} - \omega t)]$$

$$= -k^{2}\exp[i(\mathbf{k} \cdot \mathbf{r} - \omega t)]$$



#### (1931) - 3.565(3)(3) - 2529 () - 2550

## a. Plane waves and interference

For monochromatic wave (single angular frequency  $\omega$ )

Verifying plane wave solutions of the form when k=  $\omega$  /c

Since 
$$\frac{\partial^2}{\partial t^2} \exp[i(\mathbf{k}\cdot\mathbf{r} - \omega t)] = -\omega^2 \exp[i(\mathbf{k}\cdot\mathbf{r} - \omega t)],$$
  
with  $\nabla^2 \exp[i(\mathbf{k}\cdot\mathbf{r} - \omega t)] = -k^2 \exp[i(\mathbf{k}\cdot\mathbf{r} - \omega t)]$   
 $\nabla^2 \phi(x, y, z, t) - \frac{1}{c^2} \frac{\partial^2 \phi(x, y, z, t)}{\partial t^2} = 0$   
 $\mathbf{k} = \omega/c,$   
 $\mathbf{k} = \omega/c,$ 

 $\phi(x, y, z, t) = \exp[i(\mathbf{k} \cdot \mathbf{r}) - \omega t]$  is a solution for any vector direction **k** 

provided 
$$|\mathbf{k}| = \frac{\omega}{c}$$



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## a. Plane waves and interference

#### Wave interference

#### Linearity of wave equation



A plane wave solution  $\mathbf{k}_1$ 



A plane wave solution  $k_1 + k_2$ 



A plane wave solution  $\mathbf{k}_2$ 

Wave interference between 2 waves is also a solution of wave equation because of Linearity of wave equation





### What is Diffraction?

Diffraction refers to various phenomena which occur when a wave encounters an **obstacle** or a **slit**.

Interference is a phenomenon in which two waves **superpose** to form a resultant wave of greater, lower, or the same amplitude.

### Huygens' principle



Christiaan Huygens (1629-1695)

*Every unobstructed point on a wavefront will act a source of secondary spherical waves. The new wavefront is the surface tangent to all the secondary spherical waves.* 





## **b.** Diffraction

#### Waves from an aperture



By putting huygens' sources between the aperture, wave from an aperture can be modeled.

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## **b.** Diffraction

#### Waves from a point source/an aperture



As the number of point sources increase, we have converged on a good description in this model!







## b. Diffraction

### Diffraction angle

If we look far away from a set of sources,



High frequency speakers : very directional Low frequency speakers : not very directional

참고 사이트 http://yjh-phys.tistory.com/1459

#### What does the back-scattered light look like?





**Concentric circle** 



#### What does the back-scattered look like?





Zeroth order diffraction

All of the sources add up in phase to generate a phase going back

An upward or downward direction is also possible







Larger scatterers seperation gives beam closer in angle



#### Calculate the angle $\Theta$ of diffracted waves



Larger scatterers seperation gives beam closer in angle

#### Calculate the angle $\Theta$ of diffracted waves



First order diffraction

Second order diffraction

(negative) First order diffraction

Wave equation in 3 dimensions

$$\nabla^2 \phi(x, y, z, t) - \frac{1}{c^2} \frac{\partial^2 \phi(x, y, z, t)}{\partial t^2} = 0$$

Plane wave solution for monochromatic wave (single angular frequency  $\omega$ )

$$\phi(x, y, z, t) = \exp[i(\mathbf{k} \cdot \mathbf{r} - \omega t)]$$

**Diffraction angle (by an aperture)** 

$$\theta \sim \frac{\lambda}{d}$$

Angle of diffracted waves by a periodic structure

$$\sin \theta_n = \frac{n\lambda}{s}$$