

Unitary and Hermitian operators

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Unitary operators to change representations of vectors

$$|f_{old}\rangle = \begin{bmatrix} c_1 \\ c_2 \\ c_3 \\ \vdots \end{bmatrix} \quad \text{With orthonormal} \\ \text{basis function} \\ |\psi_n\rangle$$



$$|f_{new}\rangle = \begin{bmatrix} d_1 \\ d_2 \\ d_3 \\ \vdots \end{bmatrix} \quad \text{With new orthonormal} \\ \text{basis function} \\ |\phi_m\rangle$$

Coordinate
transform
matrix

$$\hat{U} = \begin{bmatrix} u_{11} & u_{12} & u_{13} & \cdots \\ u_{21} & u_{22} & u_{23} & \cdots \\ u_{31} & u_{32} & u_{33} & \cdots \\ \vdots & \vdots & \vdots & \ddots \end{bmatrix}$$

where $u_{ij} = \langle \phi_i | \psi_j \rangle$

$$|f_{new}\rangle = \hat{U} |f_{old}\rangle$$

- prove that is unitary

$$\begin{aligned} (\hat{U}^\dagger \hat{U})_{ij} &= \sum_m u_{mi}^* u_{mj} = \sum_m \langle \phi_m | \psi_i \rangle^* \langle \phi_m | \psi_j \rangle = \sum_m \langle \psi_i | \phi_m \rangle \langle \phi_m | \psi_j \rangle \\ &= \langle \psi_i | \left(\sum_m |\phi_m\rangle \langle \phi_m| \right) | \psi_j \rangle = \langle \psi_i | \hat{I} | \psi_j \rangle = \langle \psi_i | \psi_j \rangle = \delta_{ij} \\ \text{so } \hat{U}^\dagger \hat{U} &= \hat{I} \end{aligned}$$

- Hence unitary operator \hat{U} is **unitary** since its Hermitian transpose is its inverse
- Hence **any change in basis** can be implemented with a unitary operator
- We can also say that any such change in **representation to a new orthonormal basis** is a unitary transform

Unitary operators to change representations of operators



Consider a number such as $\langle g | \hat{A} | f \rangle$

With unitary \hat{U} operator to go from "old" to "new" systems

we can write
$$\begin{aligned} \langle g_{new} | \hat{A}_{new} | f_{new} \rangle &= (|g_{new}\rangle)^\dagger \hat{A}_{new} |f_{new}\rangle \\ &= (\hat{U} |g_{old}\rangle)^\dagger \hat{A}_{new} (\hat{U} |f_{old}\rangle) = \langle g_{old} | \hat{U}^\dagger \hat{A}_{new} \hat{U} | f_{old} \rangle \end{aligned}$$

Under the assumption $\langle g_{new} | \hat{A}_{new} | f_{new} \rangle = \langle g_{old} | \hat{A}_{old} | f_{old} \rangle$

Then we can derive \hat{A}_{new}

$$\hat{A}_{old} = \hat{U}^\dagger \hat{A}_{new} \hat{U}$$

$$\hat{U} \hat{A}_{old} \hat{U}^\dagger = (\hat{U} \hat{U}^\dagger) \hat{A}_{new} (\hat{U} \hat{U}^\dagger) = \hat{A}_{new}$$

In matrix terms, with

$$\hat{M} = \begin{bmatrix} M_{11} & M_{12} & M_{13} & \cdots \\ M_{21} & M_{22} & M_{23} & \cdots \\ M_{31} & M_{32} & M_{33} & \cdots \\ \vdots & \vdots & \vdots & \ddots \end{bmatrix} \quad \text{then} \quad \hat{M}^\dagger = \begin{bmatrix} M_{11}^* & M_{21}^* & M_{31}^* & \cdots \\ M_{12}^* & M_{22}^* & M_{32}^* & \cdots \\ M_{13}^* & M_{23}^* & M_{33}^* & \cdots \\ \vdots & \vdots & \vdots & \ddots \end{bmatrix}$$

A Hermitian operator is equal to its own Hermitian adjoint

then

$$M_{ij} = M_{ji}^*$$

Hence, the diagonal elements of a Hermitian operator must be real

- statements

$(\langle g | \hat{M} | f \rangle)^\dagger \equiv (\langle g | \hat{M} | f \rangle)^*$ For arbitrary $|f\rangle$ and $|g\rangle$ Since result is just a value

$$\begin{aligned} \text{So } (\langle g | \hat{M} | f \rangle)^* &\equiv (\langle g | \hat{M} | f \rangle)^\dagger = (\hat{M} | f \rangle)^\dagger (\langle g |)^\dagger = (| f \rangle)^\dagger \hat{M}^\dagger (\langle g |)^\dagger \\ &= \langle f | \hat{M}^\dagger | g \rangle \end{aligned}$$

In integral form

$$\int g^*(x) \hat{M} f(x) dx = \left[\int f^*(x) \hat{M} g(x) dx \right]^*$$

We can rewrite the right hand side using $(ab)^* = a^* b^*$

$$\int g^*(x) \hat{M} f(x) dx = \int f(x) \{ \hat{M} g(x) \}^* dx$$

Hence,

$$\int g^*(x) \hat{M} f(x) dx = \int \{ \hat{M} g(x) \}^* f(x) dx$$

Suppose $|\psi_n\rangle$ is a normalized eigenvector of the Hermitian operator M with eigenvalue μ_n

Then, by the definition

$$\hat{M}|\psi_n\rangle = \mu_n|\psi_n\rangle$$

therefore

$$\langle\psi_n|\hat{M}|\psi_n\rangle = \mu_n\langle\psi_n|\psi_n\rangle = \mu_n$$

But from the Hermiticity of \hat{M} we know

$$\langle\psi_n|\hat{M}|\psi_n\rangle = \left(\langle\psi_n|\hat{M}|\psi_n\rangle\right)^* = \mu_n^* \quad \text{and hence } \mu_n \text{ must be real}$$

- Orthogonality of eigenfunctions for different eigenvalues

Trivially $0 = \langle\psi_m|\hat{M}|\psi_n\rangle - \langle\psi_m|\hat{M}|\psi_n\rangle$

By associativity $0 = \left(\langle\psi_m|\hat{M}\right)|\psi_n\rangle - \langle\psi_m|\left(\hat{M}|\psi_n\rangle\right)$

Using $\left(\hat{A}\hat{B}\right)^\dagger = \hat{B}^\dagger\hat{A}^\dagger$ $0 = \left(\hat{M}^\dagger|\psi_m\rangle\right)^\dagger|\psi_n\rangle - \langle\psi_m|\left(\hat{M}|\psi_n\rangle\right)$

Using Hermiticity $\hat{M} = \hat{M}^\dagger$ $0 = \left(\hat{M}|\psi_m\rangle\right)^\dagger|\psi_n\rangle - \langle\psi_m|\left(\hat{M}|\psi_n\rangle\right)$

Using $\hat{M}|\psi_n\rangle = \mu_n|\psi_n\rangle$ $0 = \left(\mu_m|\psi_m\rangle\right)^\dagger|\psi_n\rangle - \langle\psi_m|\mu_n|\psi_n\rangle$

μ_m and μ_n are real numbers $0 = \mu_m\left(|\psi_m\rangle\right)^\dagger|\psi_n\rangle - \mu_n\langle\psi_m||\psi_n\rangle$

Rearranging $0 = (\mu_m - \mu_n)\langle\psi_m|\psi_n\rangle$

But μ_m and μ_n are different, so $0 = \langle\psi_m|\psi_n\rangle$ i.e., orthogonality

$$\begin{bmatrix} \ddots & & & & \\ \dots & -\frac{1}{2\delta x} & 0 & \frac{1}{2\delta x} & 0 & \dots \\ & \dots & 0 & -\frac{1}{2\delta x} & 0 & \frac{1}{2\delta x} & \dots \\ & & \dots & & \dots & & \ddots \end{bmatrix} \begin{bmatrix} \vdots \\ f(x_i - \delta x) \\ f(x_i) \\ f(x_i + \delta x) \\ f(x_i + 2\delta x) \\ \vdots \end{bmatrix} = \begin{bmatrix} \vdots \\ \frac{f(x_i + \delta x) - f(x_i - \delta x)}{2\delta x} \\ \frac{f(x_i + 2\delta x) - f(x_i)}{2\delta x} \\ \vdots \end{bmatrix} = \begin{bmatrix} \vdots \\ \left. \frac{df}{dx} \right|_{x_i} \\ \left. \frac{df}{dx} \right|_{x_i + \delta x} \\ \vdots \end{bmatrix}$$

Note this matrix is antisymmetric in reflection about the diagonal and it is not Hermitian

Indeed somewhat surprisingly d/dx is not Hermitian $\int f_n f_m' dx \neq (\int f_n' f_m dx)^*$

We can formally “operate” on the function $f(x)$
by multiplying it by the function $V(x)$
to generate another function $g(x) = V(x)f(x)$

Since $V(x)$ is performing the role of an operator
we can if we represent it as a diagonal matrix
whose diagonal elements are
the values of the function at each of the
different points

If $V(x)$ is real
then its matrix is Hermitian as required for \hat{H}